



Uptake of organic vapours and nitric acid on atmospheric freshly nucleated particles

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Abstract. Sulfuric acid, ammonia, and amines are believed to be key contributors to the initial steps in new particle formation in the atmosphere. However, other compounds such as organic compounds or nitric acid are believed to be important for further growth at larger sizes. In this study, we investigate the potential uptake of first-generation oxidation products from α -pinene (pinic and pinonic acid) and isoprene (trans- β -IEPOX, β 4-ISPOOH, and β 1-ISOPOOH), a potential highly oxidised molecule (HOM), formic acid, and nitric acid. The uptake is probed onto (SA)₁₀(base)₁₀ freshly nucleated particles (FNPs), where SA denotes sulfuric acid, and the bases are ammonia (AM), methylamine (MA), dimethylamine (DMA), or trimethylamine (TMA). The addition free energies were calculated at the ω B97X-D3BJ/6-311++G(3df,3pd)//B97-3c level of theory. We find favourable addition free energies of -8 to -10 kcal mol⁻¹ for the HOM, pinic acid, and pinonic acid on the less sterically hindered (SA)₁₀(AM)₁₀ and (SA)₁₀(MA)₁₀ FNPs. This suggests that isoprene oxidation products do not contribute to the early growth of FNPs, but the α -pinene products do, in accordance with their expected volatilities.

Calculating the second addition of a pinic acid molecule or pinonic acid molecule on the (SA)₁₀(AM)₁₀ FNPs, we find that pinic acid maintains its large addition free energy decrease due to its two carboxylic acid groups interacting with the other monomer, as well as the FNP. The pinonic-acid addition free energy drops to -3.9 kcal mol⁻¹ due to the weak interactions between the FNP and its carbonyl group and the lack of monomer–monomer interactions. Calculating the addition free energy under realistic atmospheric conditions, we find that the FNPs studied are too small (1.4 nm) to support the growth of the studied uptake monomers. We find that the accretion product pinyl diaterpenylic ester (PDPE; C₁₇H₂₆O₈) yields an addition free energy value of -17.1 kcal mol⁻¹. This suggests that PDPE can overcome the strong Kelvin effect of a 1.4 nm FNP and lead to spontaneous uptake under ambient conditions.

1 Introduction

Clouds play a significant role in shaping the global climate, and their formation begins when water condenses onto cloud condensation nuclei (CCN) (Boucher and Lohmann, 1995). Aerosols can act as CCN, delivering the required surface area for water to condense on, when they reach sizes of roughly 50–100 nm in diameter (Boucher and Lohmann, 1995). The cloud–aerosol interaction is, according to the Intergovernmental Panel on Climate Change (IPCC) Fifth Assessment Report, the largest cause of uncertainty in modern radiative forcing modelling (Canadell et al., 2021). Aerosols

form through two main pathways, namely either as primary particles, directly emitted into the atmosphere, or as secondary particles, which form through the clustering of gas vapours (Kulmala et al., 2013). Around 50 % of CCN are believed to originate from the secondary process (Boucher and Lohmann, 1995; Merikanto et al., 2009), denoted new particle formation (NPF). As an example, Zhao et al. (2024) simulated the spatial distribution of CCN at the surface level and found the fraction of CCN from NPF to be around 30 %–40 % over mainland Europe and up to ≈ 60 % over the eastern United States.

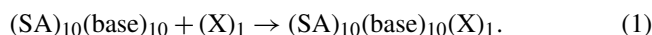
The substantial uncertainty in the climate models is linked with the unidentified mechanisms for the NPF pathways and the early growth behaviour of freshly nucleated particles (FNPs) (Canadell et al., 2021). For instance, Tröstl et al. (2016) found a factor of 2 difference in predicted CCN number concentration when altering the 1.7–3.0 nm particle growth mechanisms in their global model. It is therefore important to elucidate the initial growth mechanisms. Wang et al. (2020) found, experimentally, that vapours of ammonia (AM) and nitric acid (NA) could condense onto freshly nucleated particles (FNPs) at 278.5 K, but temperatures below 258.15 K were required for the nucleation of NA and AM. They extended this study (Wang et al., 2022) to include sulfuric acid (SA) and found that NA, SA, and AM could form particles under conditions similar to the upper free troposphere. Stolzenburg et al. (2018) studied the influence of organics in early particle growth in the CLOUD chamber and found rapid growth from organics in temperature ranges of 248.15–298.15 K.

Experimentally, the main focus has been on measuring particles larger than 2 nm in diameter, as particles below this size are challenging to measure accurately. Sub-2 nm charged-cluster compositions have been measured by mass spectrometry techniques (Jokinen et al., 2012); however, it is uncertain if the measured clusters underwent fragmentation inside the instrument or if the charging of neutral clusters changes the composition during the measurement (Zapadinsky et al., 2019; Passananti et al., 2019). This positions quantum chemical methods as a key tool for exploring the formation of small clusters and their early growth processes.

The initial clustering process is believed to primarily be driven by sulfuric acid (SA) stabilised by ammonia (AM) (Kulmala et al., 2013; Kirkby et al., 2011; Schobesberger et al., 2013; Weber et al., 1996; Elm, 2021a; Dunne et al., 2016; Kubečka et al., 2023b) or strong amines such as methylamine (MA) (Kurtén et al., 2008; Nadykto et al., 2011, 2015, 2014; Jen et al., 2014; Glasoe et al., 2015; Kubečka et al., 2023b), dimethylamine (DMA) (Kurtén et al., 2008; Loukonen et al., 2010; Nadykto et al., 2011, 2015, 2014; Jen et al., 2014; Glasoe et al., 2015; Almeida et al., 2013; Elm, 2021a; Kubečka et al., 2023b), and trimethylamine (TMA) (Elm, 2021a; Kurtén et al., 2008; Nadykto et al., 2011; Jen et al., 2014; Nadykto et al., 2015; Glasoe et al., 2015; Kubečka et al., 2023b). Nitric acid (NA) (Wang et al., 2020; Liu et al., 2021, 2018; Kumar et al., 2018; Ling et al., 2017; Wang et al., 2022; Nguyen et al., 1997; Longsworth et al., 2023; Knattrup et al., 2023; Knattrup and Elm, 2022; Bready et al., 2022; Qiao et al., 2024) and formic acid (FA) (Bready et al., 2022; Knattrup et al., 2023; Ayoubi et al., 2023; Zhang et al., 2022, 2018; Harold et al., 2022; Nadykto and Yu, 2007) have also been shown to be involved in the initial clustering. Organics have not been definitively proven to take part in the initial clustering process, but they are known to be important for the growth (Kulmala et al., 2013; Elm et al., 2020, 2023; Engsvang et al., 2023b).

While much of the focus has been on the initial clustering process (up to eight monomers) (Elm et al., 2020, 2023; Engsvang et al., 2023b), Engsvang et al. (2023a), Engsvang and Elm (2022) and Wu et al. (2023, 2024) have pushed the studies up to large particles of up to 30–60 monomers, reaching geometric diameters of up to 2 nm, where the clusters start exhibiting more particle-like properties. Wu et al. (2024) defined the cluster-to-particle transition point as the point where the free energy per monomer starts levelling off, thereby resembling “bulk” thermodynamics and where “solvated” monomers appear in the cluster structures. They found this point to be around 10 acid–base pairs for the SA–AM/MA/DMA/TMA clusters and defined the clusters at and beyond this point as FNPs. Previously, DePalma et al. (2015) studied the uptake of pinic and pinonaldehyde on a (SA)₄(AM)₄ cluster at the AM1 level of theory. However, the early growth of realistic sizes (~ 2 nm) of FNPs has yet to be studied.

In this paper, we study the uptake of NA and common organics (denoted X) on the (SA)₁₀(AM/MA/DMA/TMA)₁₀ FNPs using quantum chemical methods for the following reaction scheme:



We chose these sizes as the clusters have all reached the cluster-to-particle transition point. We investigate the first-generation oxidation products of α -pinene, isoprene (Nozière et al., 2015), and a potential highly oxidised molecule (HOM) of α -pinene (Kurtén et al., 2016) (as suggested by COSMO-RS calculations). The studied systems are displayed in Fig. 1.

2 Methodology

2.1 Computational details

The B97-3c (Brandenburg et al., 2018) and ω B97X-D3BJ//6-311++G(3df,3pd) (Najibi and Goerigk, 2018; Ditchfield et al., 1971) calculations were performed in ORCA 5.0.4 (Neese, 2012; Neese et al., 2020; Neese, 2022). The xTB 6.4.0 (Bannwarth et al., 2021) program was used for the GFN1-xTB^{re-par} calculations. GFN1-xTB^{re-par} is a parameterisation of GFN1-xTB (Grimme et al., 2017), using the methodology suggested by Knattrup et al. (2024). This specific parameterisation was done by Wu et al. (2024), where a linear combination of the binding energy and gradient errors were minimised on a set of B97-3c FNPs. Configurational sampling was performed using ABCluster 3.2 (Zhang and Dolg, 2015, 2016; Zhang, 2022) and CREST 2.12 (Grimme, 2019; Pracht et al., 2020) in noncovalent interaction mode. The entire computational workflow and subsequent data handling were performed using the JKCS 2.1 suite of programs (Kubečka et al., 2023a). All the data are freely available in the atmospheric cluster database (Kubečka et al., 2023a);

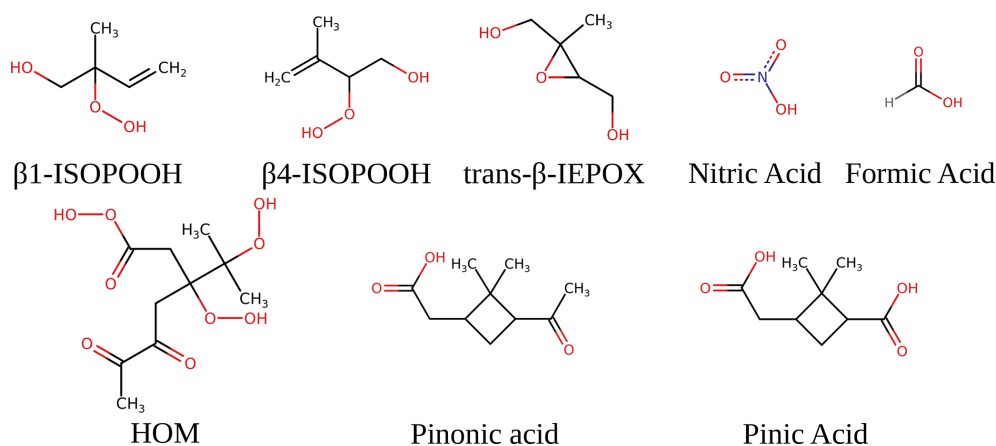


Figure 1. The molecular structure of the studied monomers. Pinonic and pinic acid are first-generation oxidation products from α -pinene. Moreover, trans- β -IEPOX, β 4-ISOPOOH, and β 1-ISOPOOH are first-generation oxidation products from isoprene.

Elm, 2019). See the “Data availability” section for more information.

2.2 Thermochemistry

We define the binding free energy ΔG_{bind} as the free energy change between the combined cluster (G_{cluster}) and the free energy (G_{monomer}) from the n monomers as follows:

$$\Delta G_{\text{bind}} = G_{\text{cluster}} - \sum_i^n G_{\text{monomer}}^i. \quad (2)$$

To get more accurate free energies, we performed a higher-level single-point correction on top of the density functional theory (DFT) geometries. This is denoted with the *sp//geo* notation, where “sp” is the single-point energy method used to calculate the electronic energy (E^{sp}), and “geo” is the method used to calculate the thermal correction to the free energy ($G_{\text{thermal}}^{\text{geo}}$) from the optimised structure and vibrational frequencies. The thermal correction term includes everything, except the electronic energy,

$$G = E^{\text{sp}} + G_{\text{thermal}}^{\text{geo}}. \quad (3)$$

The addition free energy is the binding free energy change when a monomer is added to the cluster as follows:

$$\Delta G_{\text{add}} = \Delta G_{\text{bind}}^{\text{cluster+uptake}} - \Delta G_{\text{bind}}^{\text{cluster}}. \quad (4)$$

We employ the quasi-harmonic approximation (Grimme, 2012) (as standard in ORCA), where vibrational frequencies below 100 cm^{-1} are treated as free-rotor contributions to the entropy.

$$S_{\text{rot,qh}} = \frac{1}{2}R + R \ln \left[\left(\frac{8\pi^3 I' k_B T}{h^2} \right)^{1/2} \right], \quad (5)$$

where k_B is the Boltzmann constant, T the temperature, I' the effective moment of inertia, R the gas constant, and h the Planck constant.

The output of quantum chemical programs is, for the most part, under standard conditions ($p = 1 \text{ atm}$; $T = 298.15 \text{ K}$). Changing the temperature can be done by re-evaluating the standard statistical mechanics expressions at the wanted temperature. If the binding free energy is needed under other conditions, Halonen (2022) derived the following equation for the binding free energy at the given monomer concentrations:

$$\Delta G_{\text{bind}}(p) = \Delta G_{\text{bind}}^{\text{ref}} - RT \left(1 - \frac{1}{n} \right) \sum_i^n \ln \left(\frac{p_i}{p_{\text{ref}}} \right), \quad (6)$$

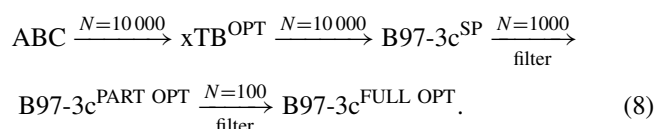
where p_i is the partial pressure (concentration) of monomer i , n is the number of monomers, and p_i is the reference pressure at which the reference binding free energy was calculated. This equation correctly satisfies self-consistency for multi-component clusters (the monomers will have a free energy of zero) and the law of mass action. Combining Eq. (6) with Eq. (4) yields the addition free energy under the given conditions, namely

$$\Delta G_{\text{add}}(p) = \Delta G_{\text{add}}^{\text{ref}} - \frac{RT}{n(n-1)} \left(\sum_i^{n-1} \ln \left(\frac{p_i}{p_{\text{ref}}} \right) + (n-1)^2 \ln \left(\frac{p_{\text{add}}}{p_{\text{ref}}} \right) \right), \quad (7)$$

where n is the total number of monomers in the largest cluster, $\Delta G_{\text{add}}^{\text{ref}}$ is the reference addition free energy, the sum runs over all monomers in the clusters before the addition, and p_{add} is the pressure (concentration) of the added monomer.

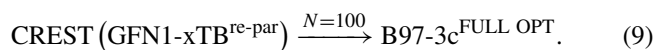
2.3 Configurational sampling workflow

We employed the configurational workflow for FNPs as suggested by Wu et al. (2023, 2024).



The workflow entails 10 parallel runs of ABCluster (Zhang and Dolg, 2015, 2016), with SN = 1280, gen = 320, and sc = 4. We used matching ionic monomers, leading to an overall electronic neutral cluster. We chose to have ionic inorganic monomers and neutral organic monomers. We chose this because SA is a stronger acid compared to the monomers. All the structures from ABCluster are subsequently optimised with GFN1-xTB^{re-par} and have a single-point energy calculated at the B97-3c level of theory. The 1000 structures that are lowest in electronic energy are then partially optimised with B97-3c for 40 iterations. The 100 structures that are lowest in electronic energy are then fully optimised, and finally, a vibrational frequency calculation is performed.

The lowest free energy is used as the input structure for CREST, using GFN1-xTB^{re-par}, as suggested by Knattrup et al. (2024), and the 100 structures that are lowest in electronic energy are then fully optimised, followed by a vibrational frequency calculation.



The larger clusters suffered from memory issues and SCF convergence problems, causing a low number of fully optimised structures using our automated workflow. We, therefore, restarted all 200 cluster calculations with increased memory and `VeryTightSCF` and redid the CREST workflow, resulting in some cluster systems having over 200 configurations. The organic monomers were sampled using CREST. The ABCluster input was generated from the lowest energy conformer, using `topgen` (Zhang, 2022; Zhang and Dolg, 2015, 2016) with charges from an MP2/6-31++G** calculation in Gaussian 16 (Gaussian et al., 2016).

2.4 Single-point refinement

Engsvang et al. (2023a) found that the B97-3c single-point energies gave quite erroneous addition free energies for FPNs compared to the DLPNO-CCSD(T₀)/aug-cc-pVTZ reference calculations. They found that $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}$ had excellent error cancellation for the addition free energies yielding errors below 1 kcal mol⁻¹. Hasan et al. (2024) extended the benchmark by including single-point calculations with $\omega\text{B97X-D3BJ}$ and the augmented def2 basis sets. However, none managed to beat the combination found by Engsvang et al. (2023a). We, therefore, performed $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}$ single-point energy calculations on the 10 lowest free energy configurations. The SCF calculations had trouble converging due to linear dependencies in the basis; therefore, we had to increase the `scfthres` parameter to 1×10^{-6} . This change slightly increases the energy (max of 0.6 kcal mol⁻¹ for the 10 acid 10 base systems); however, it is systematic for the given system size and should partly cancel out for the binding addition free energies.

It should be noted that the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}$ single-point energies slightly overbind (roughly 4 kcal mol⁻¹) compared to the DLPNO-CCSD(T₀)/aug-cc-pVTZ reference level by Engsvang et al. (2023a). However, all the energies are roughly shifted in the same manner, leading to less erroneous addition free energies.

3 Results and discussion

3.1 Addition free energies

To probe the potential for the (SA)₁₀(AM/MA/DMA/TMA)₁₀ FPNs to uptake monomers (pinic and pinonic acid, trans- β -IEPOX, β 4-ISPOOH, and β 1-ISOPOOH), HOM, FA, and NA, we calculated the addition free energies as defined in Eq. (4).

The reference FPNs, without the uptake monomers, are taken from Wu et al. (2024). They found that, for realistic atmospheric conditions ([SA] = 10⁶ molec. cm⁻³; AM = (10 ppt, 10 ppb); MA/DMA/TMA = (1, 10 ppt)), the formation of the SA-AM and SA-MA FPNs go through nucleation barriers but are stable after the barrier. Furthermore, the formation of the SA-DMA FNP is found to be entirely barrierless, but the formation of the SA-TMA FNP is unfavourable (due to steric hindrance) as the binding free energy is positive for the studied conditions.

Figure 2 presents the calculated addition free energies at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}/\text{B97-3c}$ level of theory. It should be noted that the most negative binding free energy of the FPNs follows the order of MA > DMA > AM > TMA. This trend does not align with the gas-phase acidity or basicity of the components. Instead, steric factors become increasingly important for larger sizes, thus changing the order. In the following sections, we will discuss the trends for each FNP type.

3.1.1 SA-DMA

From Fig. 2, it is seen that the (SA)₁₀(DMA)₁₀ FPNs are the least favourable for the uptake of vapours, with positive addition free energies for FA, IEPOX, and β 4-ISOPOOH and higher than -3 kcal mol⁻¹ addition free energies for the remaining monomers. This is due to the several effects. First, SA-DMA FPNs are already substantially more stable than the other FPNs due to the high basicity of DMA compared to the other bases (Wu et al., 2024). Hence, the reference system in Eq. (4) is already low in free energy, leading to higher addition free energies. Second, the two bulky methyl groups in the DMA molecules will lead to some steric hindrance when adding the organics to the FNP. Third, adding the new monomer disrupts the already favourable hydrogen bond network of the SA-DMA FPNs, either by destabilising the system or by only making it slightly more favourable. This concept is illustrated in Fig. 3 for the (SA)₁₀(DMA)₁₀(trans-

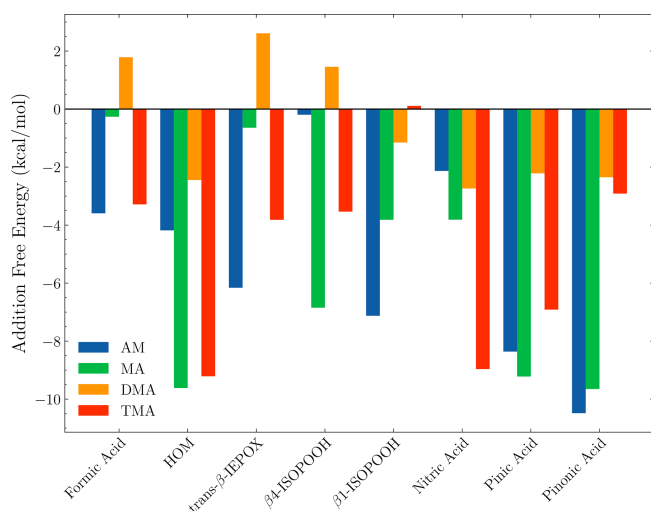


Figure 2. The addition free energy as defined in Eq. (4). The free energy is calculated using the lowest Gibbs free configuration (standard conditions) at the ω B97X-D3BJ/6-311++G(3df,3pd)//B97-3c level of theory. The x axis shows the monomer added to the $(SA)_{10}(\text{base})_{10}$ FNP, where the label defines the base.

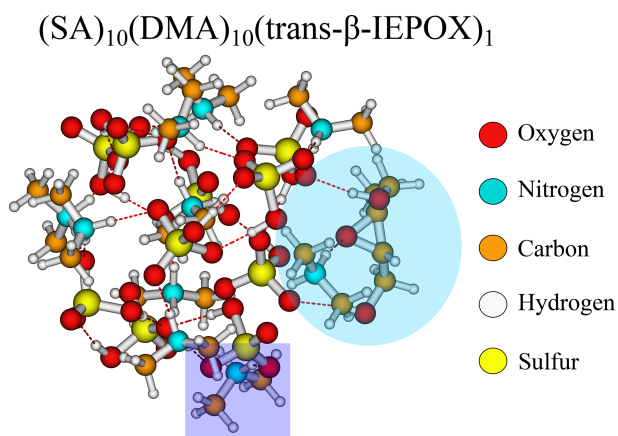


Figure 3. The $(SA)_{10}(\text{DMA})_{10}(\text{trans-}\beta\text{-IEPOX})_1$ cluster lowest in free energy (standard conditions) at the ω B97X-D3BJ/6-311++G(3df,3pd)//B97-3c level of theory.

β -IEPOX)₁ structure. The addition of the IEPOX molecule forms an epoxide–DMA bond (light-blue circle), which forces the methyl group on the DMA molecule slightly inside the cluster, leading to increased steric repulsion. This is in contrast to the preferred orientation of the DMA molecule, where both methyl groups point out of the cluster (dark-blue square).

3.1.2 SA–TMA

For the $(SA)_{10}(\text{TMA})_{10}$ FNPs, all the addition free energies are negative, except for β 1-ISOPOOH (+0.1 kcal mol^{−1}). This is mainly caused by the fact that the SA–TMA FNPs

by themselves are relatively unstable (Wu et al., 2024). Although TMA is a stronger base in terms of gas-phase basicity (Hunter and Lias, 1998) compared to DMA, its three bulky methyl groups and only a single hydrogen bond donor prevents it from obtaining a favourable hydrogen bond topological network. Hence, introducing an additional monomer extends the hydrogen bond network and lowers the steric hindrance. It is curious that there is such a large difference between the two isomers of ISOPOOH (+0.1 kcal mol^{−1} vs. −3.5 kcal mol^{−1}) as they are structurally very similar (see Fig. 4). The geometries cannot easily explain the difference as their carbon backbones point outwards from the cluster and do not interact with the “core”. This leaves the main cluster geometry as the driving parameter. Here, the $(SA)_{10}(\text{TMA})_{10}(\beta$ 4-ISOPOOH)₁ cluster seems a bit less tightly packed compared to the β 1-ISOPOOH cluster, which might equal a more favourable hydrogen bonding topological network.

Likewise, it is also curious that the bulky HOM has addition free energy in similar magnitudes (−9.2 kcal mol^{−1}) to the much smaller nitric acid (−9.0 kcal mol^{−1}) and pinic acid (−6.9 kcal mol^{−1}). Studying the structure (Fig. 5), the HOM appears to compress the entire cluster, “hovering” slightly above the surface and favouring internal hydrogen bonding as seen on top of the HOM. This means that, while the first addition is very favourable, the second addition of another HOM is less likely as this would potentially compress the cluster even further. However, the internal hydrogen bonds in the HOM could potentially be broken to act as a tether between two adjacent clusters.

3.1.3 SA–MA

The SA–MA FNPs have a notable negative addition free energy for the HOM (−9.6 kcal mol^{−1}), pinic (−9.2 kcal mol^{−1}), and pinonic acid (−9.7 kcal mol^{−1}). It is interesting that they are almost equivalent as pinic acid has two carboxylic acid groups, whereas in pinonic acid one of the carboxylic acid groups is exchanged for a carbonyl group, and we would expect carboxylic acid to bind stronger.

Studying the structures (Fig. 6), we see that they reside at the surface of the cluster, interacting with their carboxylic acid/carbonyl group in almost the same way. This suggests the clusters can rearrange to accommodate both types of interactions, and it does not matter much that one of the carboxylic acid groups has been exchanged with a carbonyl group.

Like the SA–TMA cluster, it is quite favourable to add the HOM (−9.6 kcal mol^{−1}), but, unlike the SA–TMA structure, this matches chemical intuition as the HOM almost maximises the interactions with the cluster, as seen in Fig. 7. This also fits with the interaction pattern of pinic and pinonic acid, thus explaining the magnitude of the addition free energy.

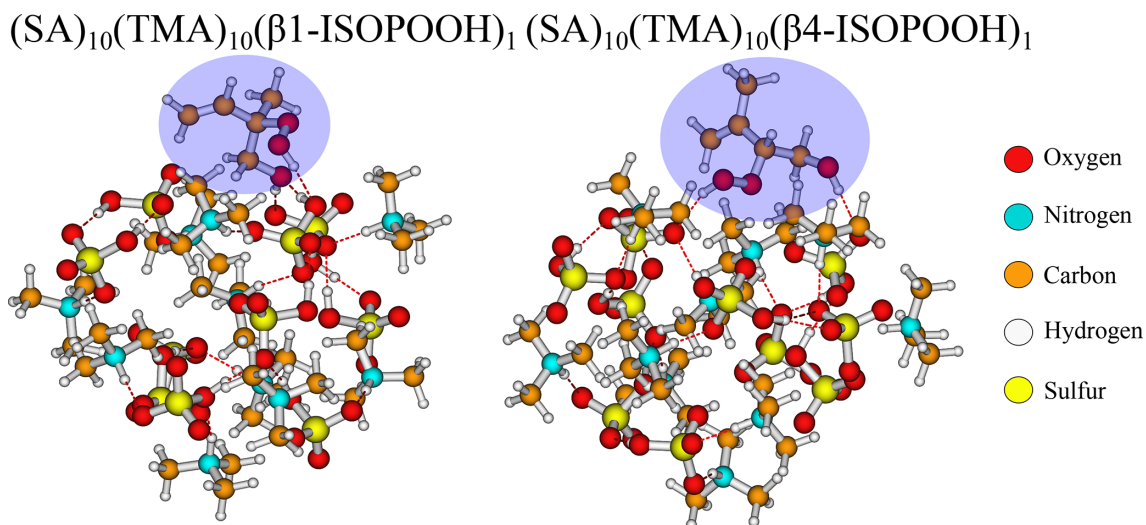


Figure 4. The $(\text{SA})_{10}(\text{TMA})_{10}(\beta\text{1-ISOPOOH})_1$ and $(\text{SA})_{10}(\text{TMA})_{10}(\beta\text{4-ISOPOOH})_1$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)//B97-3c}$ level of theory.

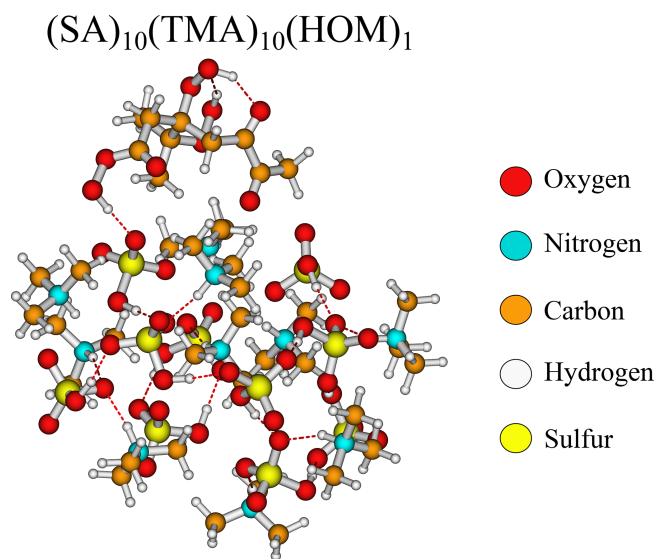


Figure 5. The $(\text{SA})_{10}(\text{TMA})_{10}(\text{HOM})_1$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)//B97-3c}$ level of theory.

3.1.4 SA–AM

The SA–AM FNPs have a similar negative addition free energy for pinic ($-8.4 \text{ kcal mol}^{-1}$) and pinonic acid ($-10.5 \text{ kcal mol}^{-1}$) to the SA–AM FNP. They also have the same structural characteristics. However, there is an inverse trend for the remaining monomers (FA, HOM, trans- β -IEPOX, $\beta\text{4-ISOPOOH}$, $\beta\text{1-ISOPOOH}$, and NA), namely if the addition free energy is substantially negative for one type of FNP, then it is small for the other type of FNP for a given monomer. This is especially noticeable for the $\beta\text{4-ISOPOOH}$

monomer, where the addition free energy is almost zero for the SA–AM FNP but $\approx -7 \text{ kcal mol}^{-1}$ for the SA–MA FNP. Structurally, it is unclear why the addition free energy differs so much as both interact with the cluster's outer layer, and neither of the “backbone” carbon chains interacts with the cluster (Fig. 8; ISOPOOH has been marked with a dark-blue box). The only major difference is that, in the SA–AM FNP, the dihedral angle between the two functional groups in ISOPOOH allows it to attach entirely to the surface, while for the SA–MA FNP, the dihedral angle between the functional groups causes it to be slightly more embedded in the cluster.

3.2 Potential for organic growth

It is clear that FNPs containing 10 acid–base pairs might have the potential to grow via organic uptake, as we observed addition free energies below -8 kcal mol^{-1} for the HOM, pinic acid, and pinonic acid. It appears that isoprene oxidation products have a lower potential for contributing to early growth compared to the α -pinene products. This is in accordance with the expected vapour pressure of the compounds.

Out of the studied organic compounds, the HOM and pinic and pinonic acid seem to be the most likely candidates for organic growth. To further investigate this, we calculated the addition free energy of a second pinic and pinonic acid molecule. We chose to study the SA–AM FNP, as it showed a large potential for the uptake of these species.

The addition free energy of the second addition is -11.8 and $-3.9 \text{ kcal mol}^{-1}$ for pinic and pinonic acid, respectively. The addition-free energy of pinic acid is even lower than the first ($-8.4 \text{ kcal mol}^{-1}$). By contrast, the addition free energy of the second pinonic acid increased dramatically from the first addition ($-10.5 \text{ kcal mol}^{-1}$). This stark contrast can eas-

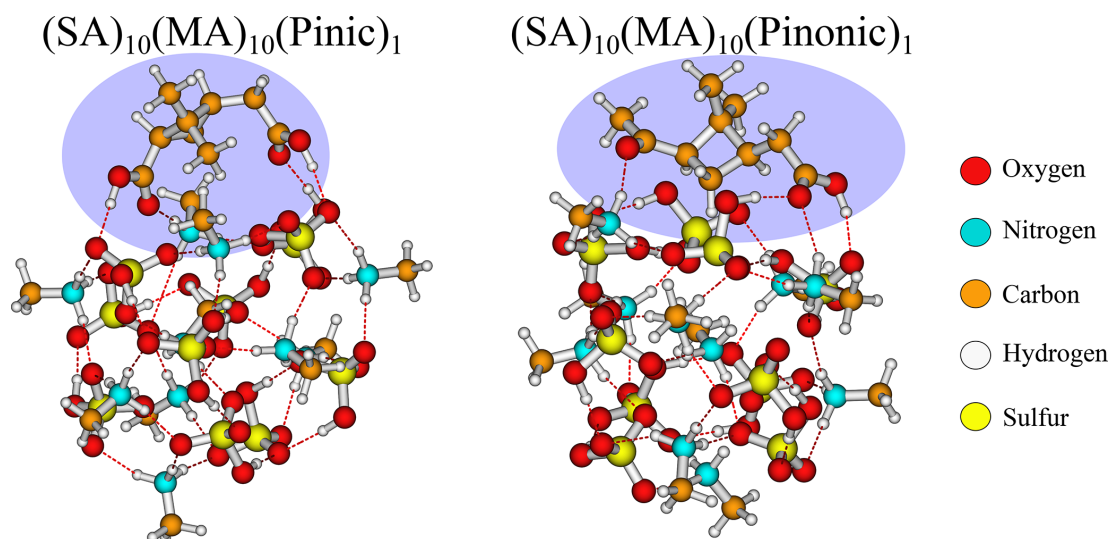


Figure 6. The $(SA)_{10}(MA)_{10}(\text{pinic})_1$ and $(SA)_{10}(MA)_{10}(\text{pinonic})_1$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}/\text{B97-3c}$ level of theory.

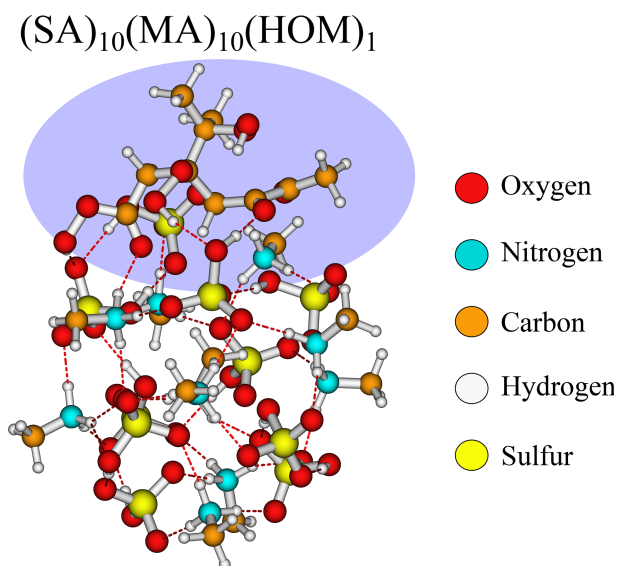


Figure 7. The $(SA)_{10}(MA)_{10}(\text{HOM})_1$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}/\text{B97-3c}$ level of theory.

ily be seen from the geometries in Fig. 9. Pinic acid, with its two carboxylic acid groups, can not only bind with the surface of the FNP but also with the other pinic acid. One could envision that with further additions of pinic acid molecules, they could link together over the entire surface, essentially leading to a core-shell structure. In contrast, pinonic acid only has one carboxylic acid group which attaches to the surface. However, it does not link the two monomers together as the carbonyl groups favours binding with the bases in the FNP. This is in agreement with the earlier findings

by Elm et al. (2017) and the cluster-of-functional-group approach by Pedersen et al. (2024), which states that carboxylic acid groups yield the most negative addition free energy for organic-enhanced atmospheric cluster formation. The fact that the organic–organic interaction is affecting the FNP growth is an intriguing result. This could indicate that the uptake of organics on small particles, and by extension the growth thermodynamics, significantly deviates from the usually employed metric of the saturation vapour pressure of the organics. Hence, uptake of organics might be as dependent on the organic–organic interactions as the organic–FNP interactions. This implies that the co-condensation of various vapours on FNPs will depend strongly on the exact functional groups in the molecules and should be further studied in the future.

3.3 Free energies under actual conditions

The free energies calculated in the previous sections are under standard conditions ($p = 1 \text{ atm}$; $T = 298.15 \text{ K}$). However, as none of the species involved is present in such high concentrations in the atmosphere, the sign of the standard free energy change alone does not determine whether the addition is spontaneous under atmospheric conditions. To explore this, we scanned the free energy dependence on temperature and monomer concentration for the addition of pinic acid onto the $(SA)_{10}(\text{AM}/\text{MA})_{10}$ FNPs. These systems were chosen because they showed favourable standard free energy changes for the first (and second, in the case of AM) addition. Two concentration regimes for the FNP components were selected based on the “clusteromics” series of papers (Elm, 2021a, b, 2022; Knattrup and Elm, 2022; Ayoubi et al., 2023). For both regimes, SA is set to $10^6 \text{ molec. cm}^{-3}$, and the concentrations of bases were AM (10 ppt; 10 ppb), MA

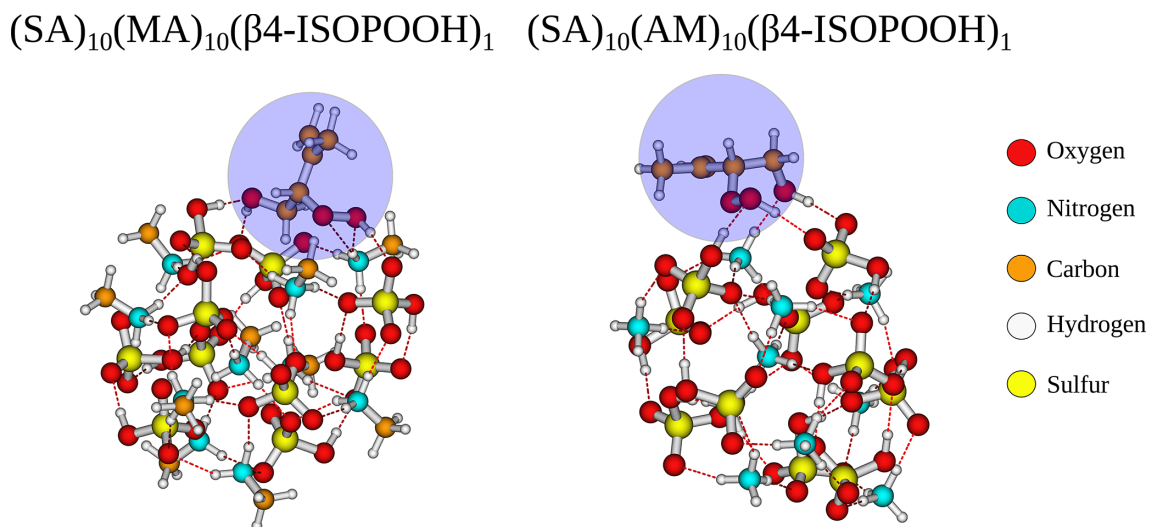


Figure 8. The $(\text{SA})_{10}(\text{MA})_{10}\beta\text{4-ISPOOH}_1$ and $(\text{SA})_{10}(\text{AM})_{10}\beta\text{4-ISPOOH}_1$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}/\text{B97-3c}$ level of theory.

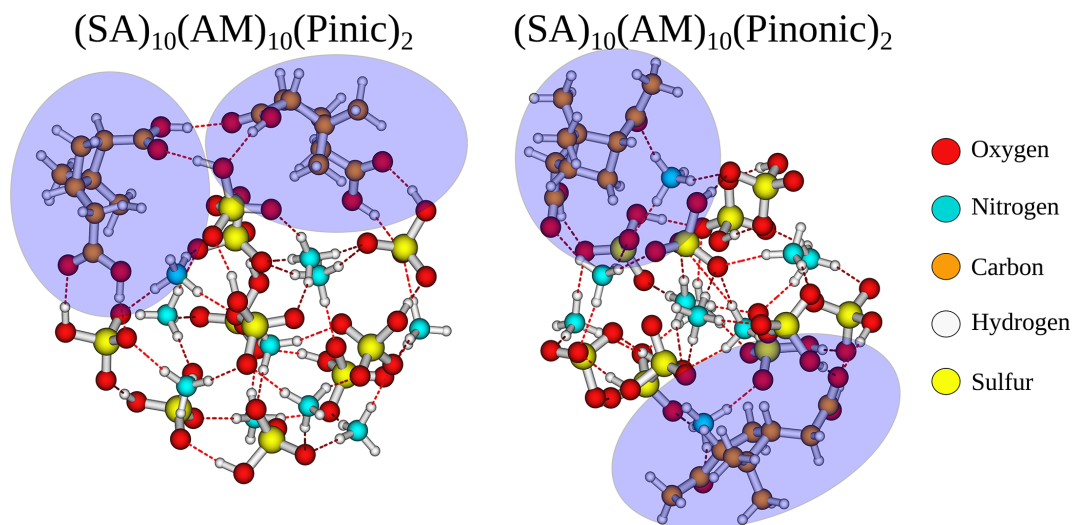


Figure 9. The $(\text{SA})_{10}(\text{AM})_{10}(\text{pinic})_2$ and $(\text{SA})_{10}(\text{AM})_{10}(\text{pinonic})_2$ cluster lowest in free energy (standard conditions) at the $\omega\text{B97X-D3BJ/6-311++G(3df,3pd)}/\text{B97-3c}$ level of theory.

(1, 100 ppt) and DMA/TMA (1, 10 ppt) for the lower and upper regimes, respectively. Ambient pinic acid concentrations were assumed to range from 0.5 to 20.5 ppt.

Figure 10 shows the maximum temperature at which the addition of pinic acid to either $(\text{SA})_{10}(\text{AM})_{10}$ or $(\text{SA})_{10}(\text{MA})_{10}$ becomes spontaneous across these concentration regimes.

In general, the primary factor determining the spontaneity of addition is the concentration of pinic acid rather than the concentration of the FNP compounds. Increasing the pinic acid concentration from 0.5 to 20 ppt raises the maximum spontaneous condensation temperature from 221 to 236 K for the addition onto the SA–AM FNP and from 231 to 246 K

for the addition onto the SA–MA FNP. Conversely, changing the base concentrations from the lower to the upper limits results in, at most, a 1 K difference. These findings suggest that pinic acid uptake is unlikely, even for the median concentration of 10 ppt, as it requires a temperature of 243 K for spontaneous addition onto the AM FNP. The addition to the MA FNP, however, is more favourable due to the direct proportionality between the addition temperature and the free energy change (Eq. 7). In the studied concentration range, the roughly $0.8 \text{ kcal mol}^{-1}$ difference in addition free energy for the AM FNP compared to MA FNP results in a 10 K increase in the condensation temperature. Hence, the results are quite sensitive to the calculated free energies.

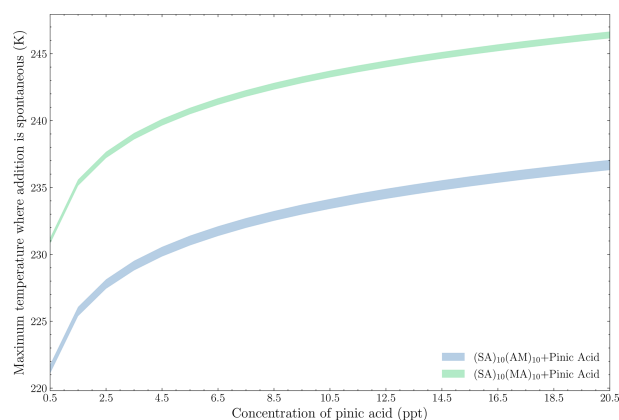


Figure 10. The maximum temperature at which the addition of pinic acid onto either (SA)₁₀(AM)₁₀ or (SA)₁₀(MA)₁₀ is spontaneous under the given conditions, ranging from the upper- to the lower-concentration regime for the bases (top of distribution = upper concentration regime; bottom of distribution = lower concentration regime).

The NA and FA uptake onto the (SA)₁₀(AM)₁₀ FNP stand out compared to the other monomers, as their higher atmospheric concentrations (ppb rather than ppt) may allow spontaneous uptake, despite their standard addition free energy exceeding -4 kcal mol^{-1} . However, even at high concentrations (20.5 ppb) in the upper-concentration regime, a spontaneous addition of FA and NA requires temperatures below 209 and 174 K, respectively. This is due to the logarithmic dependence of the free energy on concentration (from Eq. 7), which minimises the temperature increase when shifting from ppt to ppb.

These temperatures are significantly colder than those reported in CLOUD chamber experiments by Wang et al. (2020) and Wang et al. (2022), where upper-tropospheric conditions were sufficient for NA condensation. However, their experiments measured particles larger than 2 nm, whereas our SA–AM FNP is only 1.4 nm. These results suggest that monomer uptake cannot overcome the Kelvin effect for the studied FNP sizes under realistic conditions. Nonetheless, pinic acid remains the most promising candidate for spontaneous condensation onto larger FNPs under cold, but plausible, atmospheric conditions.

Another way to increase the probability of uptake is to increase the number of interacting groups in the uptake monomers. The cluster-of-functional-groups approach by Pedersen et al. (2024) states that the binding free energies of organics can, as a first approximation, be estimated by their functional groups. For uptake on the (SA)₁₀(AM)₁₀ FNP, an organic with the functional groups of pinic acid ($-8.4 \text{ kcal mol}^{-1}$) and pinonic acid ($-10.5 \text{ kcal mol}^{-1}$) would give roughly $-19 \text{ kcal mol}^{-1}$. The accretion product pinyl diaterpenylic ester (PDPE) (Kristensen et al., 2013, 2016) has three carboxylic groups and one ester group

and should approximately function as pinonic and pinic acid combined (see Fig. 11). The addition free energy for PDPE on the (SA)₁₀(AM)₁₀ FNP is $-17.1 \text{ kcal mol}^{-1}$, and the structure is depicted in Fig. 11.

When calculating the free energy under given conditions in the upper-concentration regime of FNP constituents, we find that for a concentration of 0.3 ppt for PDPE, the uptake is spontaneous at 298.15 K. At 0.03 ppt, the uptake is spontaneous at 278.15 K. This suggests that the extremely low volatile accretion products can overcome the Kelvin effect for FNPs at these sizes, but the low-volatility monomer oxidation products cannot.

4 Conclusions

We have investigated the ability of (SA)₁₀(base)₁₀ freshly nucleated particles (FNPs) with SA = sulfuric acid and base = ammonia (AM), methylamine (MA), dimethylamine (DMA), and trimethylamine (TMA) to uptake first-generation oxidation products of isoprene (trans- β -IEPOX, β 4-ISPOOH, and β 1-ISOPOOH), α -pinene (pinic and pinonic acid), a potential highly oxidised molecule (HOM), formic acid (FA), and nitric acid (NA). This was done using quantum chemical methods at the ω B97X-D3BJ/6-311++G(3df,3pd)//B97-3c level of theory.

We find that the HOM and pinic and pinonic acid can exhibit large decreases in addition free energies between -8 and $-10 \text{ kcal mol}^{-1}$, making them potential candidates for the organic growth of FNPs. This suggests that the studied isoprene oxidation products do not contribute to the early growth of FNPs, but the α -pinene products do. To further investigate this, we calculated the second addition free energy of pinic and pinonic acid onto the (SA)₁₀(AM)₁₀ FNP. We find that the pinic acid still exhibits a large addition free energy decrease of $-11.8 \text{ kcal mol}^{-1}$, but the secondary addition of pinonic acid drops to $-3.9 \text{ kcal mol}^{-1}$. The reason pinic acid maintains its high addition free energy decrease is due to its two carboxylic acid groups. The functional groups enable the pinic acid monomer to bind not only to the FNP surface but also to adjacent pinic acid. Calculating the addition free energy under realistic atmospheric conditions, we find that the FNPs studied are too small to support the growth of the studied monomer compounds, and extremely low-volatility accretion products are required for spontaneous uptake on the studied FNPs.

In the future, we imagine that molecular-dynamics-based sampling techniques, such as umbrella sampling, can give better insight into the effects of multiple minima and the partitioning of the uptake or collision simulations using molecular dynamics for realistic collision coefficients. It would also be interesting to study whether organic compounds can influence the cluster-to-particle transition points.

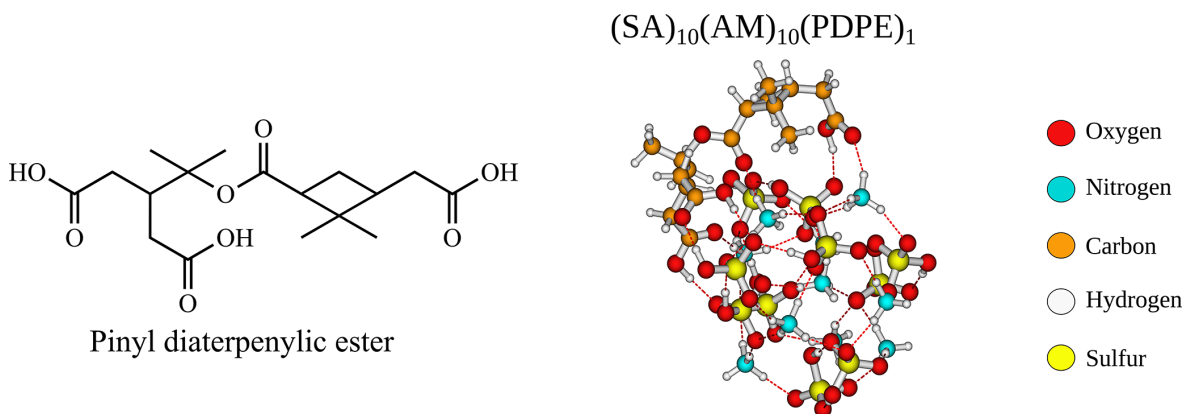


Figure 11. The structure and lowest Gibbs free configuration of the $(SA)_{10}(AM)_{10}(PDPE)_1$ FNP at the $\omega B97X-D3BJ/6-311++G(3df,3pd)//B97-3c$ level of theory.

Data availability. All the calculated structures and thermochemistry are available in the Atmospheric Cluster Database (ACDB) at <https://doi.org/10.1021/acsomega.9b00860> (Elm, 2019) and at <https://doi.org/10.1021/acsomega.3c07412> (Kubečka et al., 2023a).

Author contributions. JE: conceptualisation, project administration, funding acquisition, supervision, and resources. YK: formal analysis, investigation, writing (original draft), and visualisation. YK and JE: methodology and writing (review and editing).

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